



Washington University in St. Louis

SCHOOL OF ENGINEERING & APPLIED SCIENCE

ESE 553 - Nonlinear Dynamical Systems

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The Poincaré-Bendixson Theorem. If $\omega(x)$ is closed and bounded and does not contain an equilibrium, then $\omega(x) = \gamma$, a periodic orbit.

Corollary. If γ is a limit cycle, then there exists an open set U of initial conditions such that for $x_0 \in U$, $\omega(x_0) = \gamma$.

Corollary. If U is a bounded positively invariant open set, there is a lower bound $T_U > 0$ on the period of any periodic orbit $\gamma \subset U$.

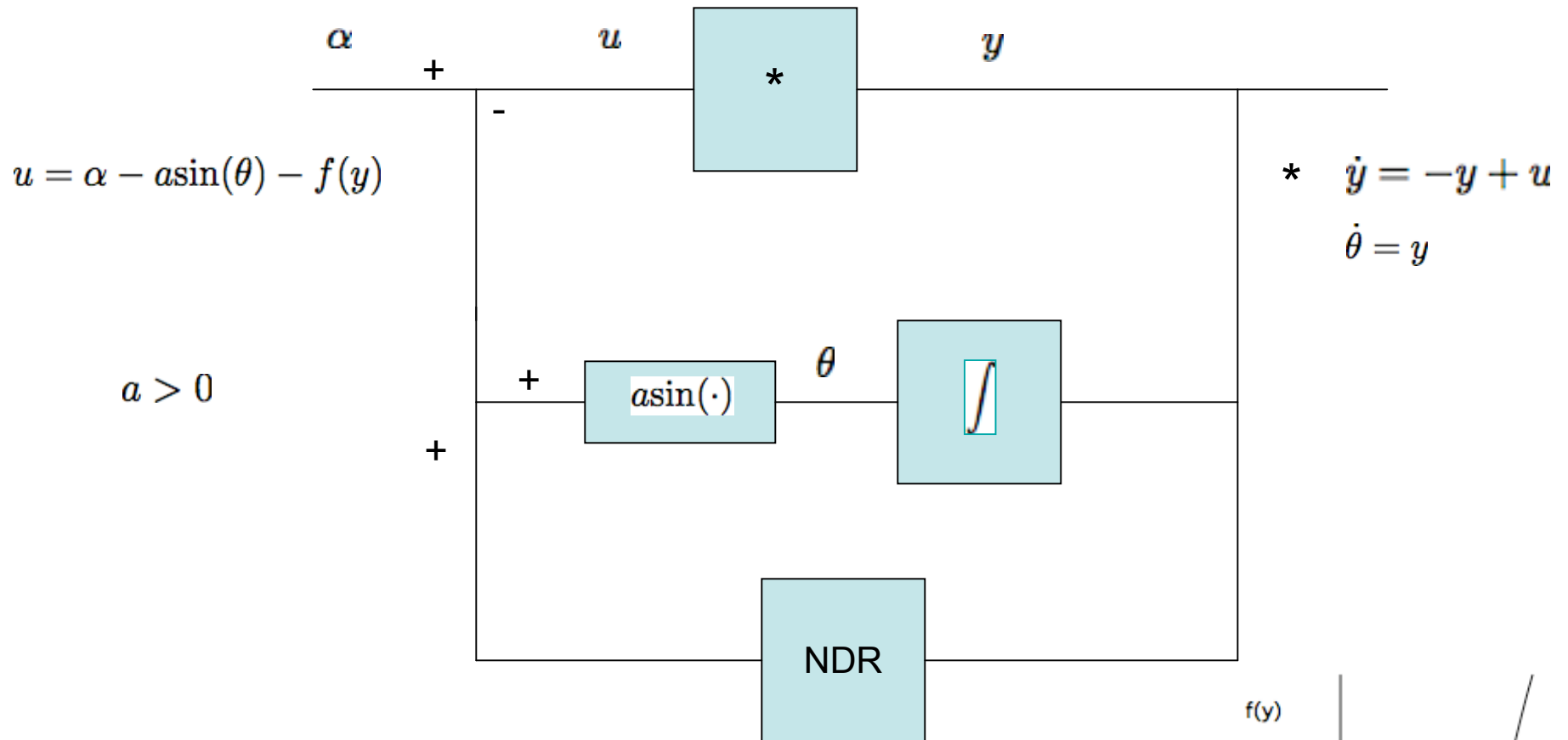
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Corollary. If γ is a periodic orbit and U is the open set consisting of the bounded interior of γ , then there exists an equilibrium $x_0 \in U$.

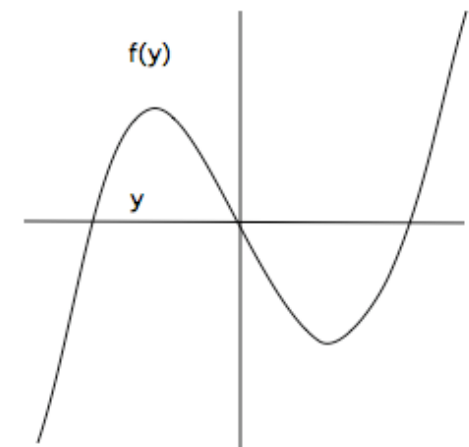
Example 1. (VCON: A Voltage Controlled Oscillator Neuron Model)



For $\theta \sim 0$ and

1. $f = 0$, classical PI controller
2. f a NDR, van der Pol.

For θ large and $f = 0$, we have the classical *Voltage Controlled Oscillator*



A Nonlinear Differential Resistor (NDR)

The van der Pol oscillator. Consider the system

$$\dot{x} = y - x^3 + \mu x \quad (1)$$

$$\dot{y} = -x. \quad (2)$$

1. If $\mu < 0$, then the origin is globally asymptotically stable.
2. For $\mu = 0$, the linearized system is a harmonic oscillator, but the origin is globally asymptotically stable.
3. For $0 < \mu \leq 1$, the system has a periodic orbit γ which is globally asymptotically stable in $\mathbb{R}^2 - \{(0, 0)\}$.

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Alternatively,

$$\dot{x} = y \quad (1)$$

$$\dot{y} = -x + \lambda(1 - x^2)y \quad (2)$$

Under the Lieńard transform, $u = x$, $v = y - \lambda(x - x^3/3)$ this becomes

$$\dot{u} = v + \lambda(u - u^3/3) \quad (1)$$

$$\dot{v} = -u \quad (2)$$

Alternatively,

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Stability of Periodic Orbits d'après Poincaré et Dulac

$$D\mathcal{P}(x_0) = \exp \left(\int_0^{t_f} \operatorname{div}(f)(\gamma(t)) dt \right)$$

Poincaré: If the integral is negative, then γ is stable.

Dulac: If there exists $\varphi : A^2 \rightarrow \mathbb{R}^+$ such that $\operatorname{div}(\varphi f) < 0$ on A^2 then there exists a unique cycle, which is asymptotically stable.

Example: (The van der Pol oscillator)

$$\left\{ \begin{array}{l} \dot{x} = y \\ \dot{y} = -x + \mu(1 - x^2)y \end{array} \right\}$$

The only equilibrium is at $(x, y) = (0, 0)$.

In fact, γ is asymptotically stable. Set

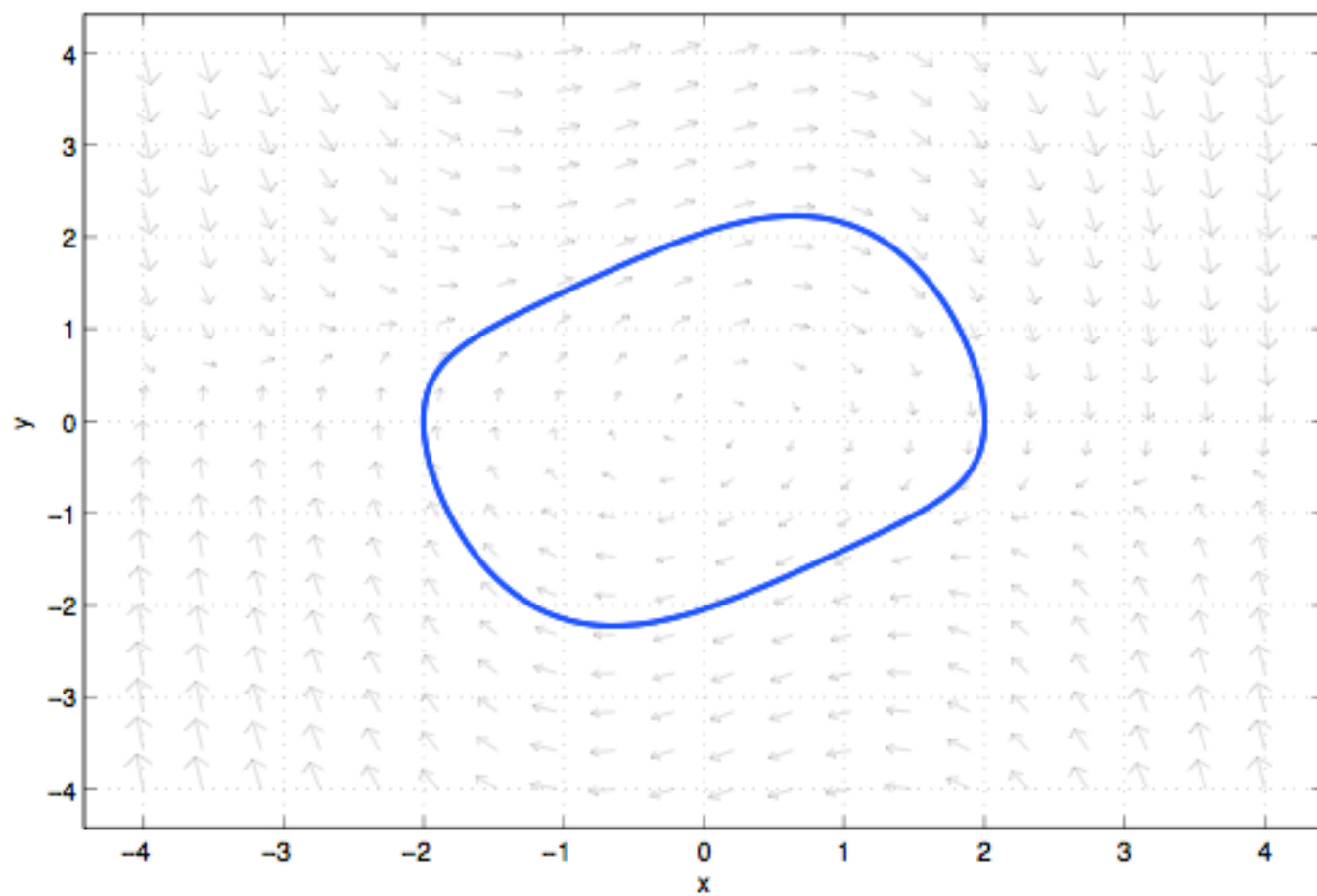
$$\phi(x, y) = (x^2 + y^2 - \rho)^{-1/2}$$

Then,

$$\frac{\partial(\phi y)}{\partial x} + \frac{\partial(-\phi x + \lambda\phi(1 - x^2)y)}{\partial y} = -\frac{\lambda(x^2 - \rho)^2}{(x^2 + y^2 - \lambda)^{3/2}} < 0$$

For $\rho = 1$ this is due to Cherkas.

$$x' = y$$
$$y' = -x + .5(1 - x^2)y$$



Steady state behavior of nonlinear systems

For B a bounded set, we define

$$\omega(B) = \{x | x = \lim_{t_n \rightarrow \infty} \phi(t_n, x_n) \text{ for } x_n \in B\}$$

Examples: The Lyapunov stable interval in the pitchfork bifurcation and the Lyapunov stable disc in the van der Pol oscillator.

The van der Pol oscillator is also undergoes a bifurcation at $\lambda = 0$. This is a complex conjugate pair of eigenvalues crossing the imaginary axis - a Hopf bifurcation.